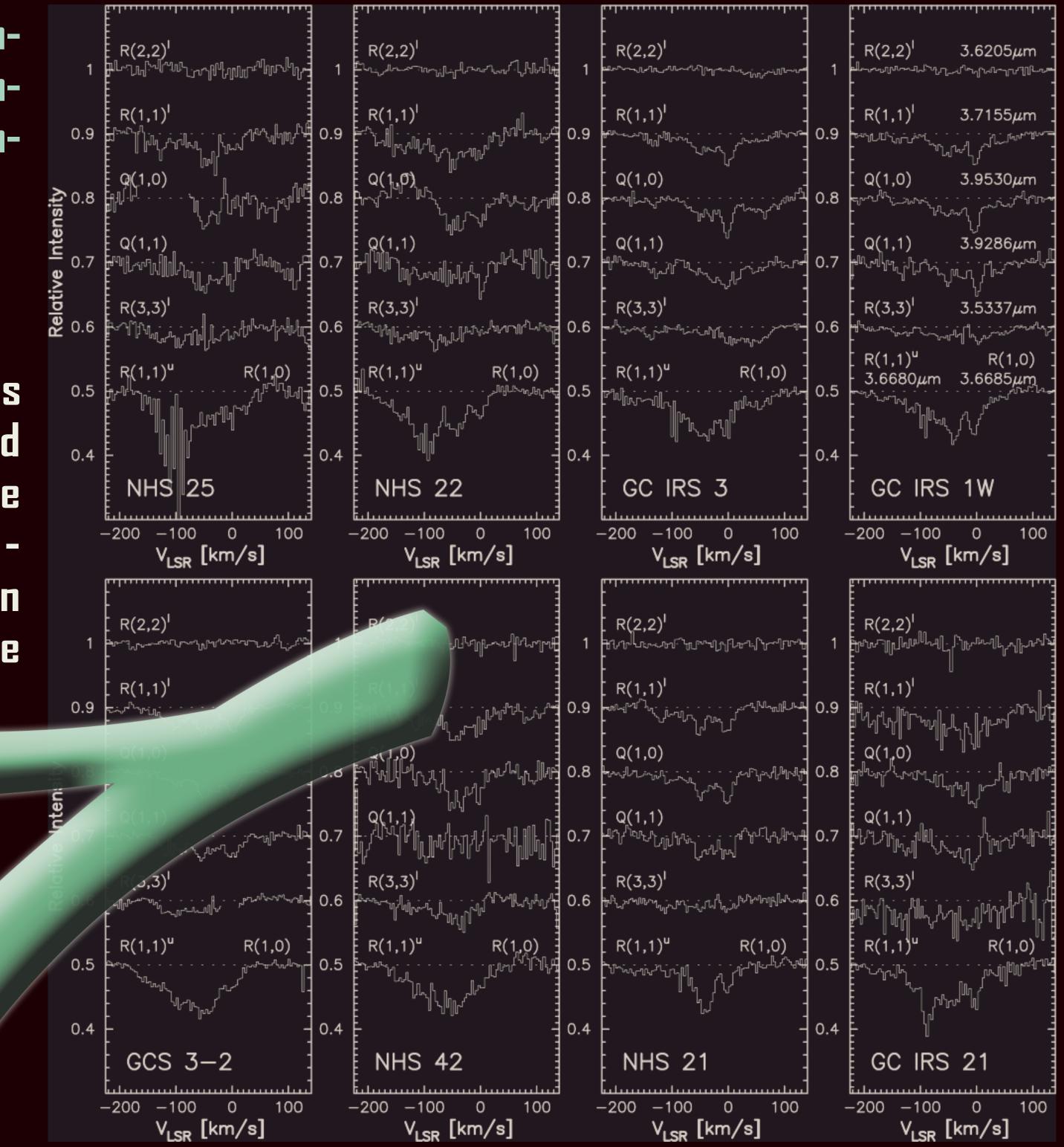
The abundance of H_3^+ in the diffuse clouds is controlled by two processes: the formation of ionization of H_2 by cosmic rays, and dissociative recombination by the electron.

In the steady state, the number density of H_3^+ is given by $\zeta n(H_2) = n(H_3^+) n_e^+ k_e^-$, where ζ and k_e^- , is the cosmic ray ionization rate and the rate constant of dissociative re- combination with the electron. The density of H_3^+ therefore product of ζ and the dimension of the cloud in



the radial dimension as $\zeta L = N(H_3^+) n_e / n(H_2) k_e$.

The column density of H_3^+ toward the Galactic center is $N(H_3^+) = (2-6)$ x 10¹⁵ cm⁻². If we take the cosmic ray ionization rate $\zeta = 2x10^{-17}$ s⁻¹ measured locally in the solar system, it gives the radial dimensions of the absorbing clouds over 1 kpc, that exceeds well beyond the boundary of the central molecular zone.

The ionization rate therefore must be as high as $\zeta = 10^{-15}$ s⁻¹ in the Galactic center (Oka et al. 2005, ApJ, 632, 882), although we still have no clues as to what is the additional source of ionization.

The relative column density of H_3^+ in (J, K) = (1,1), (2,2) and (3,3) tells that the H_3^+ in the Galactic center is in highly non-thermal rotational distribution. This makes H_3^+ a good thermometer and a dentiometer (Oka & Epp, 2004, ApJ, 613, 340). The recent detection of $\mathcal{R}(2,2)^t$ toward GC IRS 3 in the central cluster indicates that the cloud in

front of the source, M -0.02-0.07 ("50 km s⁻¹cloud"), is hot [*T* (H₂) = 450 K] and diffuse [*n*(H₂) = 200 cm⁻³] (Goto et al. 2007, in prep).

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